

**Wave equation in one spatial variable – Cauchy problem.**

*General solution of homogeneous equation*

Consider the equation

$$u_{tt} = c^2 u_{xx}, \quad x \in \mathbb{R}, \quad t > 0$$

Factorize it:

$$\left( \frac{\partial}{\partial t} - c \frac{\partial}{\partial x} \right) \left( \frac{\partial}{\partial t} + c \frac{\partial}{\partial x} \right) u = 0$$

Make the change of variables

$$\xi = x + ct, \quad \eta = x - ct$$

Then

$$\begin{aligned} \frac{\partial}{\partial x} &= \frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta}, & \frac{\partial}{\partial t} &= c \frac{\partial}{\partial \xi} - c \frac{\partial}{\partial \eta}, \\ \frac{\partial}{\partial t} - c \frac{\partial}{\partial x} &= -2c \frac{\partial}{\partial \eta}, & \frac{\partial}{\partial t} + c \frac{\partial}{\partial x} &= 2c \frac{\partial}{\partial \xi} \end{aligned}$$

The equation is transformed to the form

$$-4c^2 \frac{\partial}{\partial \xi \partial \eta} u = 0 \implies u_{\xi \eta} = 0$$

Integrate with respect to  $\eta$ :

$$u_\xi = h(\xi)$$

where  $h$  is an arbitrary function.

Integrate with respect to  $\xi$ :

$$u = \int h(\xi) d\xi + F_2(\eta) = F_1(\xi) + F_2(\eta)$$

where  $F_2$  is an arbitrary function and  $F_1(\xi) = \int h(\xi) d\xi$ .

Therefore, we can represent the general solution of the equation  $u_{tt} - c u_{xx} = 0$  in the form

$$u(x, t) = F_1(x + ct) + F_2(x - ct)$$

where  $F_1$  and  $F_2$  are two arbitrary functions.

In order to guarantee that  $u$  is a classical solution of the equation  $u_{tt} - c u_{xx} = 0$  we must assume that  $F_1$  and  $F_2$  are twice continuously differentiable.

Verification.

$$\begin{aligned} \partial_t^2 u - c^2 \partial_x^2 u &= \partial_t^2 [F_1(x + ct) + F_2(x - ct)] - c^2 \partial_x^2 [F_1(x + ct) + F_2(x - ct)] = \\ &= c^2 F_1''(x + ct) + (-c)^2 F_2''(x - ct) - c^2 F_1''(x + ct) - c^2 F_2''(x - ct) = 0. \end{aligned}$$

The term  $F_2(x - ct)$  represents the wave propagating to the right

The term  $F_1(x + ct)$  represents the wave propagating to the left

Lines

$$x - ct = C, \quad x + ct = C, \quad C \in \mathbb{R}$$

are called the *characteristics* of the wave equation

$$u_{tt} - c u_{xx} = 0.$$

*Cauchy problem for homogeneous equation*

$$\begin{aligned}u_{tt}(x, t) - c^2 u_{xx}(x, t) &= 0, \quad x \in \mathbb{R}, \quad t > 0 \\u(x, 0) = \varphi(x), \quad u_t(x, 0) &= \psi(x), \quad x \in \mathbb{R}\end{aligned}$$

We will proceed from the general solution

$$u(x, t) = F_1(x + ct) + F_2(x - ct)$$

Setting  $t = 0$  we have

$$\varphi(x) = u(x, 0) = F_1(x) + F_2(x)$$

Differentiate:

$$\varphi'(x) = F_1'(x) + F_2'(x)$$

On the other hand,

$$u_t(x, t) = c F_1'(x + ct) - c F_2'(x - ct)$$

This yields

$$\psi(x) = u_t(x, 0) = c F_1'(x) - c F_2'(x)$$

We obtain the system for  $F_1'$  and  $F_2'$ :

$$\begin{aligned}F_1'(x) + F_2'(x) &= \varphi'(x) \\ c F_1'(x) - c F_2'(x) &= \psi(x)\end{aligned}$$

Solve it:

$$\begin{aligned}F_1'(x) &= \frac{1}{2}\varphi'(x) + \frac{1}{2c}\psi(x) \\ F_2'(x) &= \frac{1}{2}\varphi'(x) - \frac{1}{2c}\psi(x)\end{aligned}$$

Integrate:

$$\begin{aligned}F_1(x) &= \frac{1}{2}\varphi(x) + \frac{1}{2c} \int_0^x \psi(\tau) d\tau + A \\ F_2(x) &= \frac{1}{2}\varphi(x) - \frac{1}{2c} \int_0^x \psi(\tau) d\tau + B\end{aligned}$$

where  $A$  and  $B$  are constants.

Taking  $x = 0$  and adding we have

$$F_1(0) + F_2(0) = \varphi(0) + A + B$$

Since  $F_1(x) + F_2(x) = \varphi(x)$ , it holds  $A + B = 0$ .

Putting the formulas of  $F_1$  and  $F_2$  into the general solution we obtain

$$u(x, t) = \frac{1}{2} [\varphi(x + ct) + \varphi(x - ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(\tau) d\tau$$

This is the *d'Alembert's formula*.

*Theorem.* Let  $\varphi \in C^2$ ,  $\psi \in C^1$ . The Cauchy problem

$$u_{tt}(x, t) - c^2 u_{xx}(x, t) = 0, \quad x \in \mathbb{R}, \quad t > 0$$

$$u(x, 0) = \varphi(x), \quad u_t(x, 0) = \psi(x), \quad x \in \mathbb{R}$$

has a unique classical solution  $u \in C^2$  given by d'Alembert's formula.

*Principle of causality.*

Initial condition at the point  $(x_0, 0)$  can spread only to that part of the  $xt$ -plane which lies between the lines with equations  $x \pm ct = x_0$  (the characteristics passing through the point  $(x_0, 0)$ ).

The sector with these boundary points is called the *domain of influence* of the point  $(x_0, 0)$ .

Let  $(x, t)$  be an arbitrary point. The values of initial conditions on the  $x$ -axis (for  $t = 0$ ) influencing the value of the solution at a point  $(x, t)$  are:

$\varphi(x - ct)$ ,  $\varphi(x + ct)$  and  
 $\psi(\tau)$  for  $\tau \in [x - ct, x + ct]$ .

The triangle  $\Delta_{xt}$  with vertices at the points  $(x - ct, 0)$ ,  $(x + ct, 0)$  and  $(x, t)$  is called the *domain of dependence* (or the *characteristic triangle*) of the point  $(x, t)$ .

*Cauchy problem for nonhomogeneous equation*

*Theorem.* Let  $\varphi \in C^2$ ,  $\psi \in C^1$ ,  $f \in C^1$ . The Cauchy problem

$$\begin{aligned} u_{tt}(x, t) - c^2 u_{xx}(x, t) &= f(x, t), \quad x \in \mathbb{R}, t > 0 \\ u(x, 0) &= \varphi(x), \quad u_t(x, 0) = \psi(x), \quad x \in \mathbb{R} \end{aligned} \quad (1)$$

has a unique classical solution  $u \in C^2$  given by the formula

$$\begin{aligned} u(x, t) &= \frac{1}{2} [\varphi(x + ct) + \varphi(x - ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(\tau) d\tau + \\ &+ \frac{1}{2c} \iint_{\Delta_{xt}} f(y, s) dy ds \end{aligned} \quad (2)$$

Here

$$\iint_{\Delta_{xt}} f(y, s) dy ds = \int_0^t \int_{x-c(t-s)}^{x+c(t-s)} f(y, s) dy ds$$



*Two methods* to derive the solution formula (2).

1. Usage of Green's theorem.

Green's theorem formulated for  $\Delta_{xt}$ : For any  $P, Q \in C^1$ ,

it holds

$$\iint_{\Delta_{xt}} [P_y - Q_s] dy ds = \int_{\partial\Delta_{xt}} P ds + Q dy$$

where  $\int_{\partial\Delta_{xt}}$  is the second kind line integral  
(counterclockwise).

Denote  $\partial\Delta_{xt} = L_0 \cup L_1 \cup L_2$  where

$L_0$  is the line segment from point  $(x - ct, 0)$  to point  $(x + ct, 0)$ ,

$L_1$  is the line segment from point  $(x + ct, 0)$  to point  $(x, t)$ ,

$L_2$  is the line segment from point  $(x, t)$  to point  $(x - ct, 0)$ .

From the differential equation, we have

$$\iint_{\Delta_{xt}} f(y, s) dy ds = \iint_{\Delta_{xt}} [u_{ss} - c^2 u_{yy}] dy ds$$

Setting  $P = -c^2 u_y$  and  $Q = -u_s$  in the Green's theorem we deduce

$$\iint_{\Delta_{xt}} [u_{ss} - c^2 u_{yy}] dy ds = \int_{L_0 \cup L_1 \cup L_2} -c^2 u_y ds - u_s dy$$

On  $L_0$ , we have  $s = 0$ ,  $ds = 0$  and  $u_s(y, 0) = \psi(y)$ , thus

$$\int_{L_0} -c^2 u_y ds - u_s dy = - \int_{x-ct}^{x+ct} \psi(y) dy$$

On  $L_1$ , we have  $y + cs = x + ct$ , thus  $dy + c ds = 0$  and

$$\begin{aligned} \int_{L_1} -c^2 u_y ds - u_s dy &= c \int_{L_1} u_y dy + u_s ds = c \int_{L_1} du = \\ &= c [u(x, t) - u(x + ct, 0)] = cu(x, t) - c\varphi(x + ct) \end{aligned}$$

On  $L_2$ , we have  $y - cs = x - ct$ , thus  $dy - c ds = 0$  and

$$\begin{aligned} \int_{L_2} -c^2 u_y ds - u_s dy &= -c \int_{L_2} u_y dy + u_s ds = -c \int_{L_2} du = \\ &= -c [u(x - ct, 0) - u(x, t)] = -c\varphi(x - ct) + cu(x, t) \end{aligned}$$

Putting pieces together, we obtain

$$\begin{aligned} \int_{L_0 \cup L_1 \cup L_2} -c^2 u_y ds - u_s dy &= \\ &= 2c u(x, t) - c[\varphi(x - ct) + \varphi(x + ct)] - \int_{x-ct}^{x+ct} \psi(y) dy \end{aligned}$$

and

$$\begin{aligned} \iint_{\Delta_{xt}} f(y, s) dy ds &= \\ &= 2c u(x, t) - c[\varphi(x - ct) + \varphi(x + ct)] - \int_{x-ct}^{x+ct} \psi(y) dy \end{aligned}$$

Expressing  $u(x, t)$  we obtain the formula

$$\begin{aligned} u(x, t) &= \frac{1}{2} [\varphi(x + ct) + \varphi(x - ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(\tau) d\tau + \\ &+ \frac{1}{2c} \iint_{\Delta_{xt}} f(y, s) dy ds \end{aligned}$$

and the formula (2) is derived.

## 2. Operator method.

Firstly, we solve the following Cauchy problem for ordinary differential equation:

$$\frac{d^2}{dt^2}v(t)+A^2v(t) = f(t), t > 0, \quad v(0) = \varphi, \quad \frac{d}{dt}v(0) = \psi \quad (3)$$

where  $A$ ,  $\varphi$  and  $\psi$  are given numbers.

The homogeneous equation is

$$\frac{d^2}{dt^2}v_H(t) + A^2v_H(t) = 0$$

The characteristic equation  $\lambda^2 + A^2 = 0$  has the solutions  $\lambda_{12} = \pm Ai$ .

The general solution of the homogeneous equation is

$$v_H(t) = C_1 \cos At + C_2 \sin At$$

We use the variation of constants to derive a particular solution of the inhomogeneous equation

$$v_P(t) = D_1(t) \cos At + D_2(t) \sin At$$

Standard system for the derivatives of the coefficients  $D_1, D_2$ :

$$\begin{aligned} D_1' \cos At + D_2' \sin At &= 0 \\ D_1'(\cos At)' + D_2'(\sin At)' &= f \end{aligned} \implies \begin{aligned} D_1' \cos At + D_2' \sin At &= 0 \\ -D_1' A \sin At + D_2' A \cos At &= f \end{aligned}$$

Determinant is

$$\text{Det} = \begin{vmatrix} \cos At & \sin At \\ -A \sin At & A \cos At \end{vmatrix} = A(\cos^2 At + \sin^2 At) = A$$

Thus

$$D_1' = \frac{\begin{vmatrix} 0 & \sin At \\ f & A \cos At \end{vmatrix}}{\text{Det}} = \frac{-f(t) \sin At}{A}$$

$$D_2' = \frac{\begin{vmatrix} \cos At & 0 \\ -A \sin At & f \end{vmatrix}}{\text{Det}} = \frac{f(t) \cos At}{A}$$

and

$$D_1(t) = -\frac{1}{A} \int_0^t \sin A\tau f(\tau) d\tau$$

$$D_2(t) = \frac{1}{A} \int_0^t \cos A\tau f(\tau) d\tau$$

The formula of  $v_P$  is

$$\begin{aligned} v_P(t) &= D_1(t) \cos At + D_2(t) \sin At = \\ &= \frac{1}{A} \int_0^t [-\sin A\tau \cos At + \cos A\tau \sin At] f(\tau) d\tau = \\ &= \frac{1}{A} \int_0^t \sin A(t - \tau) f(\tau) d\tau \end{aligned}$$

The general solution of the equation (3) is

$$v(t) = v_H(t) + v_P(t) = C_1 \cos At + C_2 \sin At + \frac{1}{A} \int_0^t \sin A(t - \tau) f(\tau) d\tau$$

Use the initial conditions  $v(0) = \varphi$ ,  $\frac{d}{dt}v(0) = \psi$ :

$$v(0) = C_1 \cos 0 + C_2 \sin 0 = C_1 \implies C_1 = \varphi$$

$$\frac{d}{dt}v(t) = -C_1 A \sin At + C_2 A \cos At + \int_0^t \cos A(t - \tau) f(\tau) d\tau$$

$$\frac{d}{dt}v(0) = -C_1 A \sin 0 + C_2 A \cos 0 = C_2 A \implies C_2 = \frac{\psi}{A}$$

The solution of the Cauchy problem (3) is

$$v(t) = \varphi \cos At + \psi \frac{1}{A} \sin At + \frac{1}{A} \int_0^t \sin A(t - s) f(s) ds$$

We write this solution in the form

$$v(t) = \frac{d}{dt}S(t)\varphi + S(t)\psi + \int_0^t S(t - s) f(s) ds$$

where

$$S(t) = \frac{1}{A} \sin At, \quad \frac{d}{dt}S(t) = \cos At$$

In the particular case  $\varphi = 0$ ,  $f \equiv 0$ , the solution is

$$v(t) = S(t)\psi \quad \text{where} \quad S(t) = \frac{1}{A} \sin At$$

Let us compare it with the d'Alembert's formula in case  $\varphi = 0$ :

$$u(x, t) = \mathcal{S}(t)\psi = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(\tau) d\tau$$

The operator  $\mathcal{S}(t)$  is called *the sine family*.

Using this analogy, we construct the solution of the Cauchy problem for the wave equation in the form

$$u(x, t) = \frac{\partial}{\partial t} \mathcal{S}(t)\varphi + \mathcal{S}(t)\psi + \int_0^t \mathcal{S}(t-s)f(x, s)ds$$

The operator  $\frac{\partial}{\partial t} \mathcal{S}(t)$  is called *the cosine family*.



Compute:

$$\begin{aligned}\frac{\partial}{\partial t} \mathcal{S}(t)\varphi &= \frac{1}{2c} \frac{\partial}{\partial t} \int_{x-ct}^{x+ct} \varphi(\tau) d\tau = \\ &= \frac{1}{2c} \frac{\partial}{\partial t} \left[ \int_0^{x+ct} \varphi(\tau) d\tau - \int_0^{x-ct} \varphi(\tau) d\tau \right] = \\ &= \frac{1}{2c} [c\varphi(x+ct) - (-c)\varphi(x-ct)] = \\ &= \frac{1}{2} [\varphi(x+ct) + \varphi(x-ct)]\end{aligned}$$

$$\begin{aligned}\int_0^t \mathcal{S}(t-s) f(x, s) ds &= \int_0^t \frac{1}{2c} \int_{x-c(t-s)}^{x+c(t-s)} f(y, s) dy ds = \\ &= \frac{1}{2c} \iint_{\Delta_{xt}} f(y, s) dy ds\end{aligned}$$

Let us put the solution together:

$$\begin{aligned} u(x, t) &= \frac{\partial}{\partial t} \mathcal{S}(t) \varphi + \mathcal{S}(t) \psi + \int_0^t \mathcal{S}(t-s) f(x, s) ds = \\ &= \frac{1}{2} [\varphi(x+ct) + \varphi(x-ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(\tau) d\tau + \\ &+ \frac{1}{2c} \iint_{\Delta_{xt}} f(y, s) dy ds. \end{aligned}$$

This is the formula (2).